

Probability and quantum foundation

J.F. Geurdes

Abstract: A classical probabilistics explanation for a typical quantum effect in Hardy's paradox is demonstrated.

1. Classical probability

The common claim is that a classical probability triple, (Ω, \mathbf{F}, P) cannot explain quantum effects. Here the sample space Ω is *any* non-empty set. The σ -field, \mathbf{F} is obtained from the set of all subsets, $\mathbf{P}(\Omega) = 2^\Omega$, of Ω . \mathbf{F} is called a σ -field [1] if, (i) $\Omega \in \mathbf{F}$, (ii) $E \in \mathbf{F} \Rightarrow E^c = (\Omega - E) \in \mathbf{F}$, (iii) $E, F, \dots \in \mathbf{F} \Rightarrow E \cup F \cup \dots \in \mathbf{F}$. The triple is completed with a probability measure P , such that, $(\forall : X \in \mathbf{F}) (0 \leq P(X) \leq 1), P(\Omega) = 1$.

2. Pre-measurement characteristics, numerals and algebra

Let us inspect the possibilities of classical probability for Hardy's paradox [2] where quantum particles like electron and positron can be measured *after* mutual annihilation. This appears to reject the possibility of pre-measurement characteristics [3].

Apart from zero, unity and two the numerals of von Neuman and of Zermelo [4] are disjoint. This fact may represent mutual exclusion of electron and positron. We have, $D_0 = C_0 = \emptyset$. Von Neuman numerals are $(n = 0, 1, 2, 3, \dots)$

$$(1) \quad C_{n+1} = \{C_0, C_1, \dots, C_n\}.$$

Hence, $C_1 = \{\emptyset\}$, $C_2 = \{\emptyset, \{\emptyset\}\}$, $C_3 = \{\emptyset, \{\emptyset\}, \{\emptyset, \{\emptyset\}\}\}$, etc.

Zermelo's system is

$$(2) \quad D_{n+1} = \{D_n\}.$$

Hence, $D_1 = \{\emptyset\}$, $D_2 = \{\{\emptyset\}\}$, $D_3 = \{\{\{\emptyset\}\}\}$, etc.

We establish mutual exclusion (annihilation) with C_3 , modeling particle 1 and D_3 modeling particle 2, for instance, as $C_3 \cap D_3 = \emptyset$.

The sample space equals $\Omega = C_3 \cup D_3$, or

$$(3) \quad \Omega = \{\emptyset, \{\emptyset\}, \{\emptyset, \{\emptyset\}\}, \{\{\emptyset\}\}\}.$$

This entails the σ -field $\mathbf{F} = \mathbf{P}(\Omega) = 2^\Omega$. Explicitly:

$$\mathbf{F} = \{\Omega, \emptyset, \{\emptyset\}, \{\emptyset, \{\emptyset\}\}, \{\emptyset, \{\emptyset, \{\emptyset\}\}\}, \{\emptyset, \{\{\emptyset\}\}\}, \{\{\emptyset\}, \{\emptyset, \{\emptyset\}\}\}, \{\{\emptyset\}, \{\{\emptyset\}\}\}\} \cup \{\{\{\emptyset, \{\emptyset\}\}, \{\{\emptyset\}\}\}, \{\{\emptyset, \{\emptyset\}\}, \{\emptyset, \{\emptyset\}\}\}, \{\{\emptyset, \{\emptyset\}\}, \{\{\emptyset\}\}\}, \{\{\emptyset\}, \{\emptyset, \{\emptyset\}\}, \{\{\emptyset\}\}\}\} \cup \{\{\emptyset, \{\emptyset, \{\emptyset\}\}, \{\{\emptyset\}\}\}, \{\{\emptyset\}\}, \{\{\emptyset, \{\emptyset\}\}\}, \{\{\{\emptyset\}\}\}\}$$

Note, $D_3, C_3 \in \mathbf{F}$. The probability measure for Ω is $P(X) = |X|/|\Omega|$, with, $|X|$ the cardinality of X , i.e. $P \sim Uniform(\Omega)$. (Ω, \mathbf{F}, P) establishes classical probability. Finally, let us introduce the 'union of a set' [5], [6] operation,

$$(4) \quad \cup Z = \{x | (\exists : y \in Z)(x \in y)\}$$

3. Application to Hardy's physics

For set $A = \{C_2\}$ and $B = D_3$ we see $A \subset C = C_3$ and $B = D = D_3$, hence, $B \subset D_3$. Obviously, $A \cap B = \emptyset$. The A and B represent disjoint parts of the electron and positron. Now, $C_2 \in A$ and that means, $\{\emptyset, \{\emptyset\}\} \in A$. Hence (eq. 4), $x = \emptyset$ and $x = \{\emptyset\}$ in C_2 giving $\cup A = C_2 = \{\emptyset, \{\emptyset\}\} \in \mathbf{F}$. Identically, $D_2 \in B = D_3$. Hence, $x = \{\emptyset\} = D_1$, such that $\cup B = D_2 = \{\{\emptyset\}\} \in \mathbf{F}$. Now, $C_2 \cap D_2 = \{\{\emptyset\}\} \Rightarrow P(C_2 \cap D_2) \neq 0$. There exist subsets of C_3 and D_3 that, after taking the union, allows for simultaneous probability $\neq 0$. Hence, classical probability can do something similar to quantum mechanics if the \cup on disjoint (sub)sets in annihilation processes is physical.

4. Conclusion

A classical probabilistics explanation for a typical quantum behavior, similar to tunneling, has been found. If \cup cannot be excluded from physics it may represent a quantum physical process and establishes a classical explanation. A possible physical picture for \cup can perhaps be associated to a 'dark' mirror-matter sector [7], [8], [9] that may arise as a consequence of the experimentally established weak interaction parity non-invariance [10], [11].

References

- [1] ROSENTHAL, J. (2006). *A first look at Rigorous Probability Theory* World Scientific, Singapore
- [2] HARDY, L. (1992). Quantum mechanics, local realistic theories and Lorentz-invariant local realistic theories *Phys. Rev. Lett.* **68** 2981-2984.
- [3] EINSTEIN, A. PODOLSKY, B. ROSEN, N. (1935). Can quantum-mechanical description of reality be considered complete? *Phys. Rev.* **47** 777-780.
- [4] JAQUETTE, D. (2002) *Philosophy of mathematics: An anthology* Blackwell Publ., Oxford UK.
- [5] RANDAL HOLMES, M. (2009). *Elementary set theory with a universal set* Volume 10 of the cahiers du centre de logique Academia Louvain-la-Neuve(Belgium).
- [6] HAJNAL, A. & HAMBURGER, P. (1999). *Set Theory* London Mathematical Society Student Texts **48**, Camb. Univ. Press Cambridge (UK).
- [7] FOOT, R. (2004). Mirror matter-type dark matter *Int. J. Mod. Phys.* **D13**, 2161-2192, arXiv:astro-ph/0407623v1.
- [8] OKUN, B. (2007). Mirror particles and mirror matter: 50 years of speculation and search *Physics-Uspekhi* **50**(4) 380-391, arXiv:hep-ph/0606202v2.
- [9] FOOT, R. (2007) Mirror dark matter arXiv:hep-ph/07062694v1.
- [10] LEE, T.D. & YANG, C.N. (1956). Question of Parity conservation in weak interaction *Phys. Rev.***104**(1), 254-258.
- [11] AMBLER, E., HAYWARD, R.W., HOPES, D.D., HUDSON, R.R., & WU, C.S. (1956). Experimental test of parity conservation in beta decay *Phys. Rev.***105**, 1413-1414.